

# Matrix Product State Description of Halperin States

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(see also arXiv:1806.06858)



# Motivations

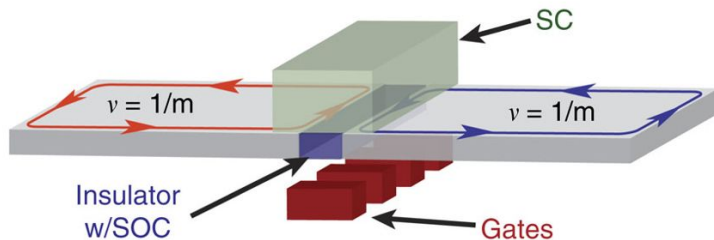
- ▶ Many FQHE experimental realizations involve an internal degree of freedom

Spin, Singlet-Triplet in QW, Graphene ...

PRB 92, 075410 (Balram et al. - 2015)

- ▶ Bilayer systems: Non-Abelian excitation at the boundary between Abelian states

Nat. Comm. 4, 1348 (Clarke et al. - 2013)



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# Motivations

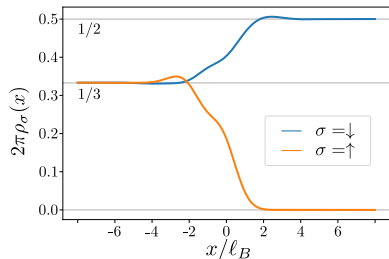
- ▶ Quasi-electronic operators for FQH states shares some of the caveats

J. Stat. Mech. 053101 (Kjäll et al. - 2018)

- ▶ Use of conformal embedding  $SU(2)_k \hookrightarrow SU(3)_k$  to create interfaces between topological orders (spinful problem - see B. Estienne's talk).

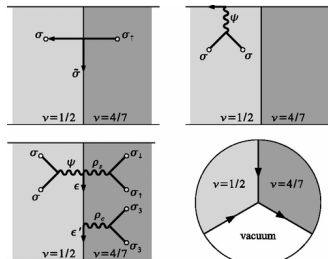
$k = 1$ : Laughlin - Halperin

V.C. et al. arXiv:1806.06858



$k = 2$ : MR - NASS

PRL 103, 076803 (Grosfeld and Schoutens 2009)



# Halperin Wavefunctions

$$\psi_{mmn} = \prod_{i < j \leq N^\downarrow} (z_i^\downarrow - z_j^\downarrow)^m \prod_{i < j \leq N^\uparrow} (z_i^\uparrow - z_j^\uparrow)^m \prod_{i,j} (z_i^\downarrow - z_j^\uparrow)^n$$

- ▶ Natural generalization of Laughlin WFs, recovered when  $N^\uparrow = 0$
- ▶ Vanishing properties encoded in the  $\mathbf{K}$ -matrix

$$\mathbf{K} = \begin{pmatrix} m & n \\ n & m \end{pmatrix}$$

PRB 46, 2290 (Wen and Zee - 1992)

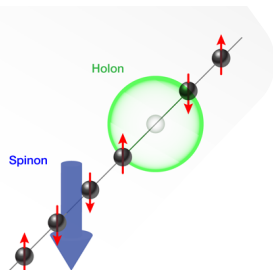
- ▶  $\psi_{mmn}$  represents the spatial WF on component  $|\uparrow \cdots \uparrow \downarrow \cdots \downarrow\rangle$ . Full WF  $|\Psi_{mmn}\rangle$  requires (anti-)symmetrization over both spin and position.
- ▶ Can be extended to any larger number of components.

# MPS Derivation 1 - CFT

The Halperin states can be written as CFT correlators.  
CFT matches the one describing the critical edge.

Nucl. Phys. B 360,2 (Moore and Read - 1991)

$$\psi_{mmn} = \langle 0 | \mathcal{O}_{\text{bkg}} \mathcal{W}^\uparrow(z_1^\uparrow) \cdots \mathcal{W}^\uparrow(z_{N\uparrow}^\uparrow) \mathcal{W}^\downarrow(z_1^\downarrow) \cdots \mathcal{W}^\downarrow(z_{N\downarrow}^\downarrow) | 0 \rangle$$



- ▶ Two-component Luttinger liquid: spin-charge separation
- ▶ Vertex operators (factorization of  $\mathbf{K} + \text{OPE}$ )

$$\mathcal{W}^\uparrow(z) =: e^{i\sqrt{\frac{m+n}{2}}\varphi^c(z) + i\sqrt{\frac{m-n}{2}}\varphi^s(z)} :$$

$$\mathcal{W}^\downarrow(z) =: e^{i\sqrt{\frac{m+n}{2}}\varphi^c(z) - i\sqrt{\frac{m-n}{2}}\varphi^s(z)} :$$

Remark: Additional  $SU(2)_1$  algebra on  $\varphi_s$  for spin singlet states (i.e.  $n = m - 1$ ).

## MPS Derivation 2 - Modes and Symmetrization

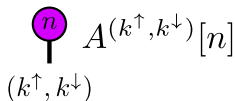
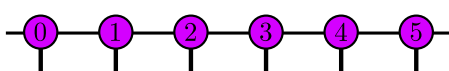
Step 1: Mode expansion  $\mathcal{W}^\sigma(z) = \sum_{n \in \mathbb{Z}} z^n W_{-n}^\sigma \Rightarrow W_{-n}^\sigma$  corresponds to an electron of spin  $\sigma$  on orbital  $n$ .

Estienne et al. arXiv:1311.2936

Step 2: Correct statistics with explicit Klein factor at the CFT level

$\chi = e^{i\pi\sqrt{2(m+n)}a_0^c} = (-1)^{Rca_0^c}$ . Numerically efficient!

$$A^{(k^\uparrow, k^\downarrow)}[n] = \frac{1}{\sqrt{k^\uparrow! k^\downarrow!}} \left(W_{-n}^\uparrow \chi\right)^{k^\uparrow} \left(W_{-n}^\downarrow\right)^{k^\downarrow}$$

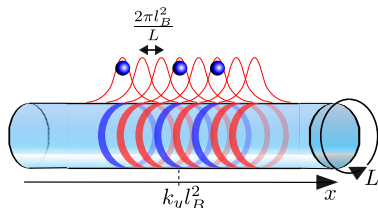


# MPS Derivation 3 - iMPS

## Cylinder geometry and iMPS

PRB 86, 245305 (Zaletel and Mong - 2012)

Spread background charge equally between orbitals and reproduce geometrical factors:



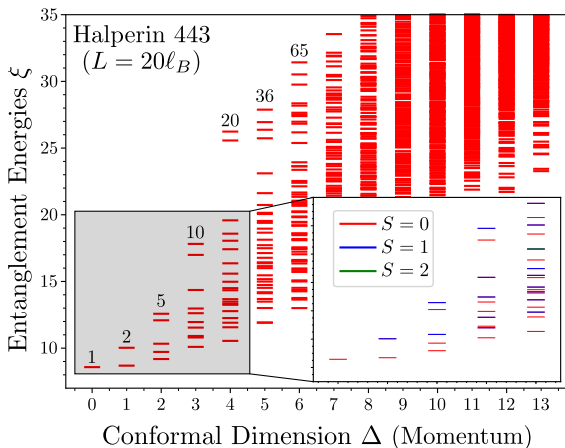
$$B(k^\uparrow, k^\downarrow) = A(k^\uparrow, k^\downarrow)[0]U$$

$$U = e^{-i\sqrt{\frac{2}{m+n}}\varphi_0^c} e^{-\left(\frac{2\pi}{L}\right)^2 L_0}$$

$L_0$  is the Hamiltonian on a circle, states becomes irrelevant with increasing energy. Natural truncation of infinite matrices *wrt* the conformal dimension (as in Truncated Conformal Space Approach of RG).

# Numerical Implementation and Entanglement Entropy

Another motivation for this truncation: it preserves the exact CFT structure.



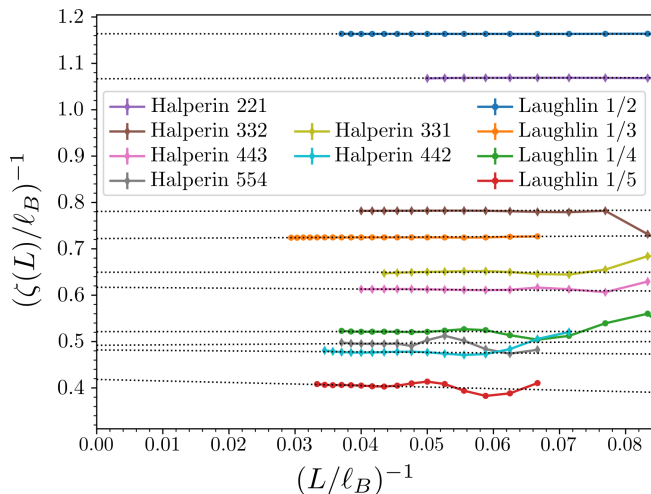
- ▶ Counting of two-component free boson 1-2-5-10-20-...
- ▶ Exact SU(2) symmetry (Ward identities)  
Even at finite truncation!
- ▶ Exact structure requires more states than usual DMRG procedure



# Numerical Results: Correlation Lengths

*Gapped?* Yes, it may describe a Hall liquid.

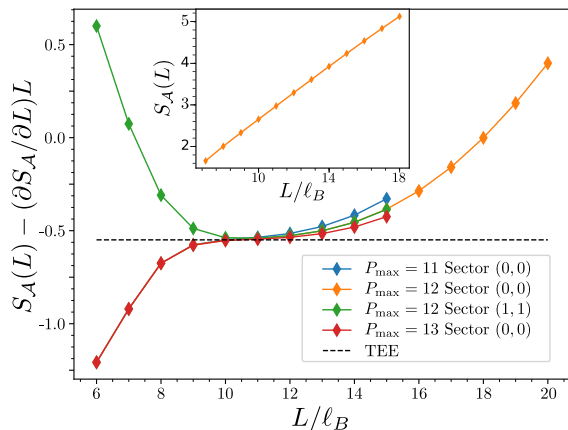
Decay of any generic correlator, necessary for braiding/adiabatic arguments!



# Numerical Results: Topological Entanglement Entropy

A fully developed topological order? Extract the correction to the area law  $S_A(L) = \alpha L - \gamma$  which characterizes the topological order.

PRL 96, 110404 (2006) ; PRL 96, 110405 (2006)



- ▶ We extract the TEE  $\gamma$  from the derivative  $S_A - L\partial_L S_A = -\gamma$
- ▶ Extracted  $\gamma \simeq 0.545(5)$  agrees with theory  $\log \sqrt{\det \mathbf{K}} = 0.549\dots$
- ▶ Abelian: Same  $\gamma$  for all topo sectors

# Conclusion

- ▶ We derived an MPS to describe the Halperin states and all their zero energy excitations.
- ▶ Solved symmetrization issue (requirement for multicomponent continuous MPS!) and reached an infinite MPS form by spreading the background charge.
- ▶ Imposed a global  $SU(2)$  symmetry by a careful choice of truncation.
- ▶ Numerical characterization of the phase topological content, and evaluation of the correlation lengths.

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# Perspectives

- ▶ Composite fermion WF and Quantum Hall hierarchies. Symmetrization relies on the same Klein factors, and quasi-electronic operators derived  
J. Stat. Mech. 053101 (Kjäll et al. -2018)
- ▶ Microscopic characterization of interfaces. Transition between equivalent  $\mathbf{K}$ -matrices?  
Cano et al., PRB 89, 115116
- ▶ Emergence of Moore-Read from polarization of a  $1/2+1/2$  bilayer. Transition between the two (inhomogeneous polarization): Majorana trapped at the interface?  
PRB 93, 085115 (Liu et al. - 2016)
- ▶ Topologically distinct boundary conditions: interface Laughlin  $1/2-1/4$  (parity of  $e/4$  quasiholes number at the transition)?

**Thank you for your attention!**

# Backup 1: Correlation Lengths

